

Microwave Theory and Techniques
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Module - 02
Lecture - 06
Waveguides - I: Parallel Plane Waveguides

Hello my name is Rajbala, I am pursuing PhD under the guidance of Professor Girish Kumar. I am also one of the TA's for this course, and I will take few lectures on waveguides. The outline to cover this topic will be.

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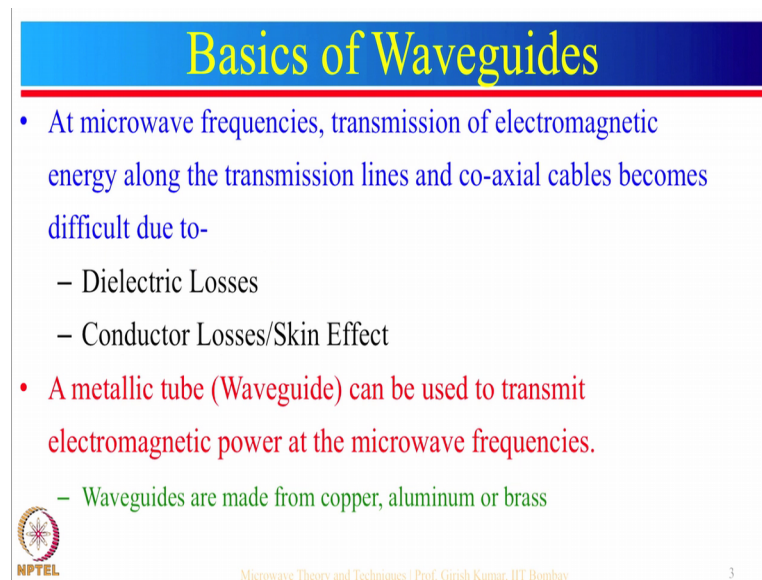
Outline

- Fundamentals of Waveguide
- Parallel Plane Waveguide
- Rectangular Waveguide
- Applications

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First I will start fundamentals of waveguides, then I will discuss parallel plane waveguide, after that I will discuss rectangular waveguide, and then I will discuss the applications of waveguides.

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Basics of Waveguides

- At microwave frequencies, transmission of electromagnetic energy along the transmission lines and co-axial cables becomes difficult due to-
 - Dielectric Losses
 - Conductor Losses/Skin Effect
- A metallic tube (Waveguide) can be used to transmit electromagnetic power at the microwave frequencies.
 - Waveguides are made from copper, aluminum or brass

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So, let us begin with basics of waveguides. The waveguides as the name suggest that it is a structure which can be used to guide electromagnetic waves along it. There are other structures also which can guide electromagnetic waves such has transmission lines and co axial cables. But, there are some differences between these structures and waveguides. And the first difference is that at microwave frequency is the transmission lines and the co-axial cables, become in efficient due to dielectric losses and conductor losses or we can say due to skin effect. Whereas, waveguides can be used at microwave frequency is and they can provide larger bandwidth with lower attenuation or lower losses.

And the second difference is that, though transmission lines can work from DC or a 0 frequency to a certain high frequency that means, it acts as a low pass filter; whereas, wave guides can be operated above a certain frequency called as cutoff frequency; so it acts as a high pass filter. And the third difference is that the transmission line can support only TEM mode of propagation whereas, wave guides can support many field configurations called as modes which will discuss later. So, these are the differences between transmission lines and waveguides.

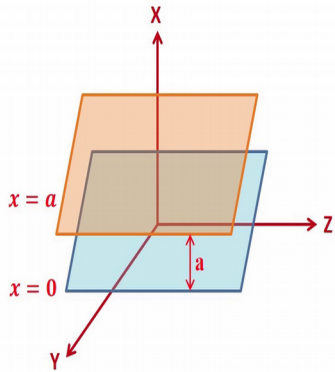
Generally, waveguides are made from high conductivity metals such has copper, aluminium, brass etcetera. And these waveguides can be of any shape, but in general the rectangular and circular waveguides are used. So, to analyze rectangular waveguide first we will develop some concept for a simple wave guiding a structure that is parallel plane

waveguide and after that we will extend those concepts to rectangular wave guide. So, let us begin with parallel plane waveguide.


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Parallel Plane Waveguide

- The plates are infinite in one direction (y- direction) and separated by a distance 'a' in x-direction.
- The direction of propagation is z-direction.



The diagram illustrates a parallel plane waveguide. It consists of two parallel plates, one at $x = a$ and one at $x = 0$, separated by a distance a in the x -direction. The plates are infinite in the y -direction. The z -axis is the direction of propagation, and the y -axis is perpendicular to the plates. The plates are shown in a 3D perspective, with the top plate in orange and the bottom plate in blue.

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In parallel plane waveguide there are two metallic plates separated by a distance a in X direction and these plates are extended to infinity in Y direction and the direction of a propagation of electromagnetic wave is taken in positive Z direction.

Now, let us see how electromagnetic wave propagate in this parallel plane waveguide. In general to analyze any problem in electromagnetics we need to solve Maxwell's equations.

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Maxwell's Equations

$$\begin{aligned} \nabla \cdot D &= \rho \\ \nabla \times E &= -\frac{\partial B}{\partial t} \quad \text{Or} \quad \nabla \times E = -j\omega\mu H \\ \nabla \cdot B &= 0 \\ \nabla \times H &= J + \frac{\partial D}{\partial t} \quad \text{Or} \quad \nabla \times H = -j\omega\epsilon E \end{aligned}$$

Longitudinal Fields: E_z, H_z

Transverse Fields in terms of longitudinal Fields:

$$\begin{aligned} H_x &= -\frac{\gamma}{h^2} \frac{\partial H_z}{\partial x} + \frac{j\omega\epsilon}{h^2} \frac{\partial E_z}{\partial y} & E_x &= -\frac{\gamma}{h^2} \frac{\partial E_z}{\partial x} - \frac{j\omega\mu}{h^2} \frac{\partial H_z}{\partial y} & \gamma &= \alpha + j\beta \\ H_y &= -\frac{\gamma}{h^2} \frac{\partial H_z}{\partial y} - \frac{j\omega\epsilon}{h^2} \frac{\partial E_z}{\partial x} & E_y &= -\frac{\gamma}{h^2} \frac{\partial E_z}{\partial y} + \frac{j\omega\mu}{h^2} \frac{\partial H_z}{\partial x} & k &= \omega\sqrt{\mu\epsilon} \\ & & & & h^2 &= \gamma^2 + k^2 \end{aligned}$$


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So, what are Maxwell's equations? There are 4 Maxwell's equations which can be written either in differential form or in integral form. The difference between these two forms is that the differential form, establish relationship between the field and the source but they cannot be used at media interphases where medium properties changes abruptly.

In those situations integral form of Maxwell's equations can be used and they will establish relationship between the fields in two different mediums they are called as boundary conditions. So, these are the 4 Maxwell's equations in differential form or they are also called as Maxwell's equation in point form.

The first Maxwell's equation is $\nabla \cdot D$ is equal to ρ which comes from Gauss law and it states that the electric flux leave in from a volume is proportional to the charge enclosed. The second equation is $\nabla \times E$ is equal to $-\frac{\partial B}{\partial t}$ which comes from the faradays law of induction, and it states that the voltage induced is proportional to the rate of change of magnetic flux. And the third equation is $\nabla \cdot B$ is equal to 0 which comes from Gauss law for magnetism and it states that total flux leaving a closed surface is 0. And the fourth and last equation is $\nabla \times H$ is equal to $J + \frac{\partial D}{\partial t}$, where j is conduction current density and $\frac{\partial D}{\partial t}$ is displacement current density.

So, these are 4 Maxwell's equations. However, for waveguides this J term in the last equation will be 0 as we assume that the wave guide is filled with the source free loss

less dielectric material. So, this J will be 0 and the last equation will reduced to del cross H is equal to dou D by dou t.

In general in waveguides we consider the direction of propagation in Z direction and the fields present in the direction of propagation are called as longitudinal fields. So, the longitudinal fields will be E z and H z and the fields which are perpendicular or transverse to the direction of propagation are called as transverse fields. So, transverse fields will be H x, H y, E x and E y.


We can derive these transverse field in terms of longitudinal fields E z and H z by solving these two Maxwell's equations del cross E is equal to minus dou B by dau t and del cross H is equal to dou D by dou t. We can consider these two equation and frequency domain which is del cross E is equal to minus j omega mu H and del cross H is equal to minus j omega epsilon E.

So, by you solving these two equation. We can get the transverse fields H x, H y, E x, E y these are the transverse fields. And in these transverse fields so constant gamma is a propagation constant which is equal to alpha plus j beta where alpha is attenuation constant and beta is phase constant and this k is equal to omega route mu epsilon and h square is equal to gamma square plus k square.

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Wave Equations

$\nabla^2 E + \omega^2 \mu \epsilon E = 0$	$\nabla^2 H + \omega^2 \mu \epsilon H = 0$
$\frac{\partial^2 E_x}{\partial x^2} + \frac{\partial^2 E_x}{\partial y^2} + \frac{\partial^2 E_x}{\partial z^2} = -\omega^2 \mu \epsilon E_x$	$\frac{\partial^2 H_x}{\partial x^2} + \frac{\partial^2 H_x}{\partial y^2} + \frac{\partial^2 H_x}{\partial z^2} = -\omega^2 \mu \epsilon H_x$
$\frac{\partial^2 E_y}{\partial x^2} + \frac{\partial^2 E_y}{\partial y^2} + \frac{\partial^2 E_y}{\partial z^2} = -\omega^2 \mu \epsilon E_y$	$\frac{\partial^2 H_y}{\partial x^2} + \frac{\partial^2 H_y}{\partial y^2} + \frac{\partial^2 H_y}{\partial z^2} = -\omega^2 \mu \epsilon H_y$
$\frac{\partial^2 E_z}{\partial x^2} + \frac{\partial^2 E_z}{\partial y^2} + \frac{\partial^2 E_z}{\partial z^2} = -\omega^2 \mu \epsilon E_z$	$\frac{\partial^2 H_z}{\partial x^2} + \frac{\partial^2 H_z}{\partial y^2} + \frac{\partial^2 H_z}{\partial z^2} = -\omega^2 \mu \epsilon H_z$



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Helmholtz Equations

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To derive a transverse fields we need to know the longitudinal fields, let us say, so let us see how longitudinal fields E_z and H_z can be found these fields E_z and H_z can be found using wave equation. And the wave equation for electric field is $\nabla^2 E + \omega^2 \mu \epsilon E = 0$ and it is $\nabla^2 H + \omega^2 \mu H = 0$ for magnetic field.

Since electric field and magnetic field has 3 components E_x, E_y, E_z and H_x, H_y and H_z . So, there will be six a scalar equations and these equations are also called as Helmholtz equation. So, to find E_z and H_z we need to solve this equation in the direction of propagation. So, we will solve these two equation $\frac{\partial^2 E_z}{\partial x^2} + \frac{\partial^2 E_z}{\partial z^2} = -\omega^2 \mu \epsilon E_z$ by $\frac{\partial^2 E_z}{\partial x^2} + \frac{\partial^2 E_z}{\partial y^2} + \frac{\partial^2 E_z}{\partial z^2} = -\omega^2 \mu \epsilon E_z$.

So, we will solve this equation and we can get E_z . Similarly if we solve this equation for magnetic field then we can get magnetic field in Z direction. Now, let us see how these longitudinal fields derived using these wave equations in parallel plane wave guide.

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Fields in Parallel Plane Waveguide

$$\frac{\partial}{\partial y} = 0$$


$$\frac{\partial^2 E_z}{\partial x^2} + \frac{\partial^2 E_z}{\partial z^2} = -\omega^2 \mu \epsilon E_z ; E_z(x, z) = X(x) Z(z)$$

$$XZ'' + XZ'' = -\omega^2 \mu \epsilon XZ \Rightarrow \frac{X''}{X} + \frac{Z''}{Z} = -k^2$$

$$\Rightarrow -k_x^2 + \gamma^2 = -k^2 \quad \text{or} \quad -k_x^2 - \beta^2 = -k^2$$

$$\frac{X''}{X} = -k_x^2 \Rightarrow X'' + k_x^2 X = 0 \Rightarrow X = C_1 \cos(k_x x) + C_2 \sin(k_x x)$$

$$\frac{Z''}{Z} = \gamma^2 \text{ or } -\beta^2 \Rightarrow Z'' - \gamma^2 Z = 0 \Rightarrow Z = C_3 e^{\gamma z} + C_4 e^{-\gamma z}$$


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So, in parallel plane waveguide as we discuss that the plates are $x \rightarrow \pm \infty$ in the direction y , so the field will be constant in y direction and that will not be any change in the fields. So, we can take $\frac{\partial}{\partial y} = 0$. By putting this thing in the wave equation we will get $\frac{\partial^2 E_z}{\partial x^2} + \frac{\partial^2 E_z}{\partial z^2} = -\omega^2 \mu \epsilon E_z$.

Since E_z is varying along x and z only and it is constant along y . So, we can consider E_z is equal to function of x into function of z . Now, put this function in this equation, so we will get Z into double derivative of X plus X into double derivative of Z is equal to minus $\omega^2 \mu \epsilon XZ$. Now, divide this whole equation with X into Z . Then we will get double derivative of X divided by X plus double derivative of Z divided by Z is equal to minus k^2 , where k^2 is equal to $\omega^2 \mu \epsilon$ this is a constant, where k is a wave number.

Since this k^2 is a constant and these variables are independent variables. So, these two terms should be a constant. So, we can take double derivative of X divided by X is equal to minus k^2 and double derivative of Z divided by Z is equal to minus β^2 or we can replace β with γ . If we take $\alpha = 0$ in γ then this will become minus β^2 . So, this negative sign for these constant is taken for the propagating wave.

And, if we take positive sign then the solution of this differential equation will be exponential function. Now, equate this minus k^2 to double derivative of X divided by X then we will get this differential equation double derivative of X plus $k^2 X$ is equal to 0.

The solution of this type of differential equation is $C_1 \cos$ function plus $C_2 \sin$ function. So, X will be $C_1 \cos kx$ plus $C_2 \sin kx$. Similarly we can find Z and this will be $C_3 e^{\gamma Z}$ plus $C_4 e^{-\gamma Z}$. So, this is how we find X and Z . Now, we will put these functions X and Z in E_z . And then we can get the longitudinal field E_z which will be $C_1 \cos kx$ plus $C_2 \sin kx$ into $C_3 e^{\gamma Z}$ plus $C_4 e^{-\gamma Z}$.

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Fields in Parallel Plane Waveguide


$$E_z(x, z) = (C_1 \cos(k_x x) + C_2 \sin(k_x x))(C_3 e^{\gamma z} + C_4 e^{-\gamma z})$$

Wave Propagation: Along +z direction $\Rightarrow C_3 = 0$

$$E_z(x, z) = (A_1 \cos(k_x x) + A_2 \sin(k_x x))e^{-\gamma z}$$

$$H_z(x, z) = (B_1 \cos(k_x x) + B_2 \sin(k_x x))e^{-\gamma z}$$

$H_x = -\frac{\gamma}{h^2} \frac{\partial H_z}{\partial x}$	$E_x = -\frac{\gamma}{h^2} \frac{\partial E_z}{\partial x}$	$\gamma = \alpha + j\beta$
		$k = \omega\sqrt{\mu\epsilon}$
		$h^2 = \gamma^2 + k^2$
		$= k_x^2$



$H_y = -\frac{j\omega\epsilon}{h^2} \frac{\partial E_z}{\partial x}$
 $E_y = \frac{j\omega\mu}{h^2} \frac{\partial H_z}{\partial x}$

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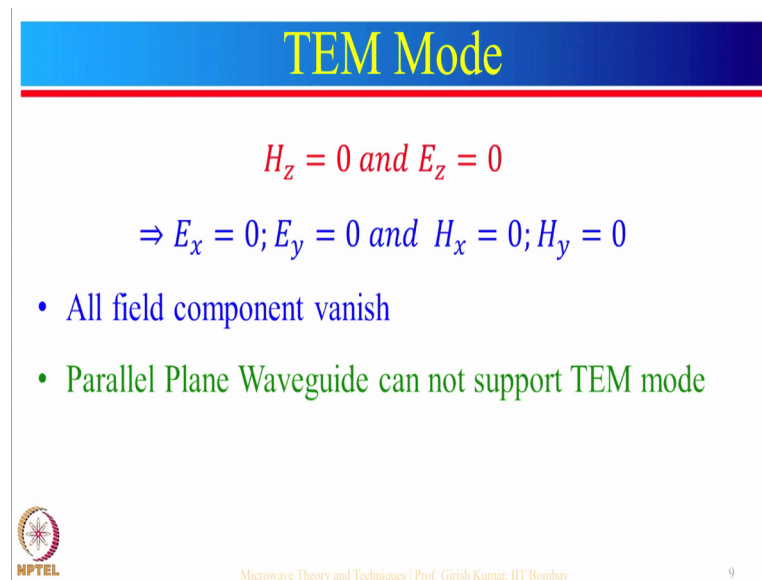
In this equation the term $E e^{\gamma z}$ represents the wave propagating in negative Z direction and the term $E e^{-\gamma z}$ represents the wave propagating in positive Z direction.

Since, we have considered the direction of propagation of electromagnetic wave in positive Z direction. So, this term should be 0 so C_3 will be 0. Now, put C_3 equal to 0 here and then E_z reduces to $C_1 \cos k_x x + C_2 \sin k_x x$ into $E e^{-\gamma z}$ or we can write it as $A_1 \cos k_x x + A_2 \sin k_x x$ into $E e^{-\gamma z}$. So, this is how longitudinal electric field is derived using wave equation, similarly we can derive magnetic field also. So, that will be $B_1 \cos k_x x + B_2 \sin k_x x$ into $E e^{-\gamma z}$.

Now, by using these two equations we can derive the transverse fields H_x , H_y , E_x and E_y . So, the equations for these transverse fields will be like this. And you can verify these by putting $\frac{\partial}{\partial y} = 0$ in the equations we have seen earlier, this is how all the field components are derived in a parallel plane waveguide.

Now, let us see how different modes propagate in a parallel plane waveguide. First let us see the TEM mode of propagation.


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TEM Mode

$$H_z = 0 \text{ and } E_z = 0$$
$$\Rightarrow E_x = 0; E_y = 0 \text{ and } H_x = 0; H_y = 0$$

- All field component vanish
- Parallel Plane Waveguide can not support TEM mode

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TEM mode is a transverse electric and magnetic mode in which electric field. And magnetic fields are transverse to the direction of propagation or there is no electric, and magnetic field in the direction of propagation that means, H_z is equal to 0 and E_z is equal to 0.

Now, by putting these values in the transverse field equations we can find the transverse fields by putting H_z equal to 0 and E_z equal to 0 we will get E_x is equal to 0, E_y is equal to 0 and H_x is equal to 0 and H_y is equal to 0. It means all the field components are 0 in a parallel plane waveguide in TEM mode is there is no field then propagation of electromagnetic wave will not take place or we can say parallel plane waveguide cannot support TEM mode of propagation.

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TM Mode

$H_z = 0; E_z \neq 0$

General Solution:


$$E_z(x, z) = (A_1 \cos(k_x x) + A_2 \sin(k_x x)) e^{-\gamma z}$$

Boundary Conditions:

(i) At $x = 0, E_z = 0 \Rightarrow A_1 = 0 \Rightarrow E_z(x, z) = E_0 \sin(k_x x) e^{-\gamma z}$

(ii) At $x = a, E_z = 0 \Rightarrow \sin(k_x a) = 0 \Rightarrow k_x a = m\pi \Rightarrow k_x = m\pi/a$

$$\Rightarrow E_z(x, z) = E_0 \sin\left(\frac{m\pi}{a} x\right) e^{-\gamma z}$$



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Now, let us see the next mode this is TM mode transverse magnetic mode in which magnetic field is transverse to the direction of propagation. And there is no magnetic field in the direction of propagation. So, H_z is equal to 0 and E_z is not equal to 0. And the general solution for electric field in the longitudinal direction will be E_z equal to $A_1 \cos k_x x$ plus $A_2 \sin k_x x$ into E raise to minus γZ as we discuss earlier.

Now, we need to apply boundary conditions on the electric field. So, what are the boundary conditions? The answer is that the tangential component of the electric field should be 0 at the conducting boundary in the parallel plane waveguide the conducting plates are at x is equal to 0 and x is equal to a and the tangential field is E_z . So, at x is equal to 0 and x is equal to a either should be 0.

So, put x is equal to 0 here we will get $A_1 \cos 0$ plus $A_2 \sin 0$ $\sin 0$ is 0, $\cos 0$ is 1. So, we will get $A_1 E$ raise to minus γZ . Now, equate this two 0 then we will get A_1 equal to 0. Now, put A_1 equal to 0 in this equation then E_z reduces to $E_0 \sin k_x x$ into E raise to minus γZ .

Now, apply the second boundary condition which is at x equal to a , E_z equal to 0. So, if you put x equal to a here then it will become $E_0 \sin k_x a$ E raise to minus γZ . And if you equate $A_2 0$ then $E_0 \sin k_x a$ cannot be 0 this cannot become 0. So, \sin function has to be 0. So, $\sin k_x a$ is 0 which implies $k_x a$ should be multiple of π . So, $k_x a$ is equal to $m\pi$ where m is an integer. From here k_x is $m\pi/a$.

Now, put this value of k_x in this equation. So, the longitudinal electric field will be E_z is equal to $E_0 \sin \frac{m\pi}{a} x e^{-\gamma z}$. So, this is how the longitudinal electric field is derived.

Now, let us see the propagating and non propagating modes in TM propagation.

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Propagating and Non-propagating TM Modes

$$H_z = 0; E_z(x, z) = E_0 \sin\left(\frac{m\pi}{a}x\right) e^{-\gamma z}$$

$$E_x = -\frac{\gamma}{h^2} \frac{m\pi}{a} E_0 \cos\left(\frac{m\pi}{a}x\right) e^{-\gamma z}; \quad E_y = 0$$

$$H_y = -\frac{j\omega\epsilon}{h^2} \frac{m\pi}{a} E_0 \cos\left(\frac{m\pi}{a}x\right) e^{-\gamma z}; \quad H_x = 0$$


Non-propagating modes:

- $TM_0: H_z = 0; E_z = 0; E_x = 0; E_y = 0; H_x = 0; H_y = 0$

Same as TEM mode

Propagating modes:

- $TM_m; m \geq 1$


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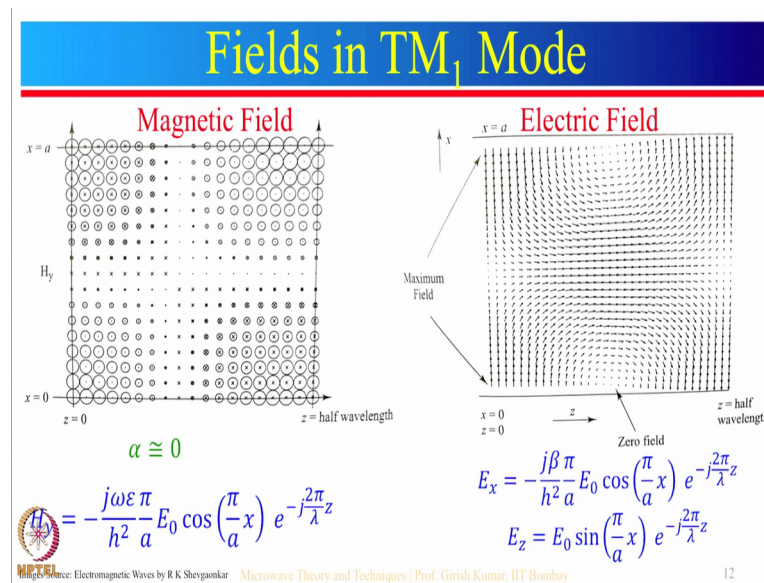
So, we have H_z is equal to 0 E_z equal to $E_0 \sin \frac{m\pi}{a} x e^{-\gamma z}$ by using these to we can find E_x , E_y , H_y , H_x . So, E_x is this thing E_y is 0, H_y is this and H_x is 0.

Now, if we put m equal to 0 in this field equations then that will become TM_0 mode, H_z is already 0. Now, put m equal to 0 in E_z then it will become $\sin 0$ and $\sin 0$ is 0. So, E_z is also 0. Then put m equal to 0 here, so 0 into something will be 0. So, E_x is 0, E_y is already 0 and H_y if we put m equal to 0 then H_y will become 0 and H_x is already 0. It means all the field components are 0. So there will not be any wave propagation.

So, this mode is same as the TEM mode because E_z and H_z both are 0 in this mode. Now, if we put m equal to 1, a more than 1 then E_z , E_x and H_y will not be 0. So, there will be propagation of electric magnetic wave in these modes. So, examples of propagating modes are TM_1 , TM_2 , TM_3 like this.

Now, let us see how field varies in these propagating modes.

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So, taken example of TM 1 mode and let us see how field varies and TM 1 mode in transverse magnetic modes, whereas only one component of magnetic field which is in y direction whereas, there are two components of electric fields which are E_x and E_z all of these 3 components are constant in the direction y and they are varying along x and Z only. So we will see though variation of the fields in XZ plane only.

So, let us first see though variation of magnetic field in XZ plane. So, in this is the X axis this is the Z direction and y direction is normal to this plane outward. And the dots in the circle represents the positive y direction and cross in the circle represents the negative y direction and the size of the circle represent the amplitude of the fields. So, larger the size of the circle higher will be the amplitude of the fields.

Now, let us see variation of magnetic field along x direction. So, H_y is varying as cos function along x direction, at x is equal to 0 this will be cos 0 which is maximum. So, at x is equal to 0 H_y will be maximum and if we put x equal to a here then this will be cos pi which is minus 1. So, at x equal to a there will be a maximum in opposite direction and if we put x equal to a by 2 then this will be cos pi by 2 which is equal to 0. So, H_y will have 0 at x is equal to a by 2.

So, the magnetic field is maximum at x is equal to 0 and x equal to a and it is 0 at x is equal to a by 2. As you can see from this figure it is maximum here maximum here 0 here. Now, let us see the variation of magnetic field in Z direction. So, it is varying as e

raise to minus $j 2 \pi$ by λ into Z in Z direction. And this function is periodic function with period 2π and it will be maximum at Z is equal to 0, Z is equal to λ by 2 and Z is equal to multiple of λ by 2. And it will be minimum at Z equal to λ by 4 Z is equal to 3λ by 4 and odd multiple of λ by 4.

So, H_y will be maximum at Z is equal to 0 Z is equal to λ by 2, Z is equal to multiple of λ by 2 and it will be 0 at Z is equal to λ by 4 3λ by 4 and odd multiple of λ by 4. So, this is how magnetic field varies in TM 1 mode. So, there is a half sinusoidal variation of magnetic field in the X direction.

Now, let us see the variation of electric field in XZ plane. So, in this the direction of arrow shows the direction of field and the length of line represents the amplitude of the fields. So, larger the length of the line the higher will be the amplitude of the field. Since, we have two components of electric field so the resultant electric field will be vector sum of E_x and E_z .

Now, let us see the variation of these two in X direction. So, E_x varies as cos function along x and E_z varies as sin function along x that means, these two components are in phase quadrature along x direction or we can say when E_x is maximum then E_z will be 0 and when E_x is 0 then E_z will be maximum. As you can see from this also so at x is equal to 0 E_x is maximum and there is no field component in Z direction, so E_z is 0.

Similarly, at x is equal to a E_x is maximum and there is no field component in Z direction, so E_z is 0. So, this is how electric field varies along X direction and in Z direction these two fields are staggered by a length of λ by 4 because of this factor there will be a 90 degree phase delay. So, when E_x is maximum then E_z will be 0, along Z direction and when E_x is 0, then E_z will be maximum along Z direction.

You can see from this figure also. So, at Z is equal to 0 E_x is maximum and E_z is 0 and at Z is equal to λ by 4, E_z is maximum and E_x is 0 and at Z is equal to λ by 2 E_x is maximum and E_z is 0. So, this is how electric and magnetic field vary in exact plane and TM 1 mode. There is half sinusoidal variation of the fields in the x direction or in the direction in which the wave is confined. So, this is all about the TM 1 mode.

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TE Mode

$$E_z = 0; H_z \neq 0$$


General Solution:

$$H_z(x, z) = (B_1 \cos(k_x x) + B_2 \sin(k_x x)) e^{-\gamma z}$$

Boundary Conditions: $E_y = \frac{j\omega\mu}{h^2} \frac{\partial H_z}{\partial x}$

(i) At $x = 0$, $E_y = 0$ or $\frac{\partial H_z}{\partial x} = 0 \Rightarrow B_2 = 0 \Rightarrow H_z(x, z) = H_0 \cos(k_x x) e^{-\gamma z}$

(ii) At $x = a$, $E_y = 0$ or $\frac{\partial H_z}{\partial x} = 0 \Rightarrow \sin(k_x a) = 0 \Rightarrow k_x a = m\pi \Rightarrow k_x = m\pi/a$

$$\Rightarrow H_z(x, z) = H_0 \cos\left(\frac{m\pi}{a} x\right) e^{-\gamma z}$$


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Now, let us move on to the next mode which is TE mode transverse electric mode in which the electric field is transverse to the direction of propagation or there is no electric field in the direction of propagation or we can say E_z is equal to 0 and H_z is not equal to 0. The general solution for magnetic field in longitudinal directional will be H_z is equal to $B_1 \cos k_x x$ into x plus $B_2 \sin k_x x$ into x into E raise to minus γZ which we derived earlier.

Now, we need to apply boundary conditions for this and the boundary conditions are the tangential component of the electric field should be 0 at the conducting boundary. So, conducting boundary is at x is equal to 0 and x is equal to a and the tangential fields are E_y and E_z . E_z is already 0 so we need to make E_y equal to 0 and to make E_y equal to 0 we need to find E_y first. So, we can find E_y in terms of H_z like this. So, E_y will be $j\omega\mu$ by h^2 into $\frac{\partial H_z}{\partial x}$. So, if E_y equal to 0 then $\frac{\partial H_z}{\partial x}$ will be 0. So, put x equal to 0 and make $\frac{\partial H_z}{\partial x}$ equal to 0 then we will get B_2 equal to 0. And the magnetic fields reduces to H_z is equal to $H_0 \cos k_x x$ into x E raise to minus γZ .

Now, apply the second boundary condition which is at x is equal to a , E_y is equal to 0 or at x is equal to a $\frac{\partial H_z}{\partial x}$ is equal to 0. By applying this condition we will get $\sin k_x x$ into x is equal to 0 which means $k_x x$ is equal to $m\pi$ by a . Now, put this value in

this equation. So, we will get the longitudinal magnetic field H_z is equal to $H_0 \cos \frac{m\pi}{a} x e^{-\gamma z}$.

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Propagating and Non-propagating TE Modes


$$E_z = 0; H_z = H_0 \cos\left(\frac{m\pi}{a}x\right) e^{-\gamma z} \Rightarrow \begin{cases} E_x = 0; E_y = -\frac{j\omega\mu}{h^2} \frac{m\pi}{a} H_0 \sin\left(\frac{m\pi}{a}x\right) e^{-\gamma z} \\ H_x = \frac{\gamma}{h^2} \frac{m\pi}{a} H_0 \sin\left(\frac{m\pi}{a}x\right) e^{-\gamma z}; H_y = 0 \end{cases}$$

Non-propagating modes:

- $TE_0: E_z = 0; H_z \neq 0; E_x = 0; E_y = 0; H_x = 0; H_y = 0$

Propagating modes:

- $TE_m; m \geq 1$


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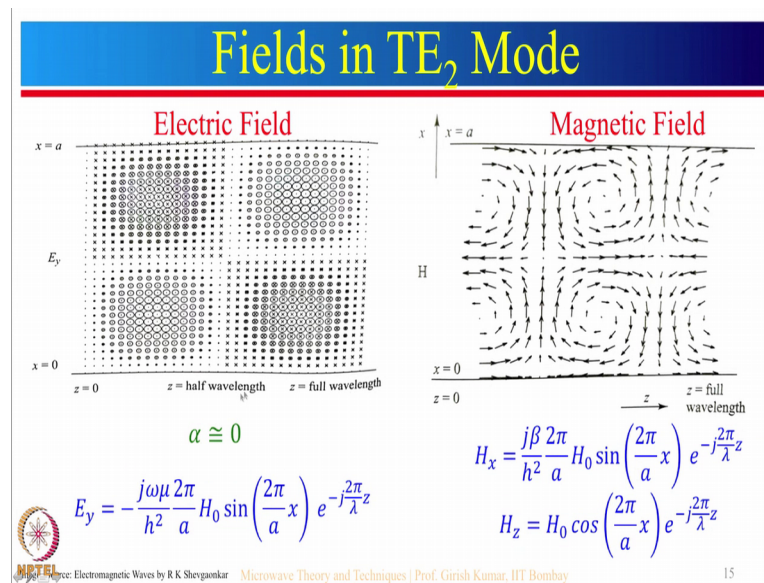
Now, let us see propagating and non propagating modes in TE propagation. So, we have E_z is equal to 0 and H_z is equal to this $H_0 \cos \frac{m\pi}{a} x e^{-\gamma z}$. From here we can find the transverse fields. So, E_x is equal to 0, E_y is equal to this, H_x is equal to this and H_y is equal to 0.

Now, if we put m is equal to 0 in these field equations then this will become TE_0 mode. In this E_z is already 0 and by putting m is equal to 0 H_z will be $H_0 e^{-\gamma z}$ which is not equal to 0. E_x is already 0, E_y will be 0 into something 0 H_x will be 0 into something 0 and H_y is already 0. So, all the transverse field components are 0. So, there will not be any wave propagation in TE_0 mode so for TE propagation non propagating mode is TE_0 .

Now, if we put m equal to 1 or greater than 1 then this H_z , E_y and H_x will not be 0. So, there will be wave propagation in TE_m mode. So, the propagating modes in TE propagation are TE_m mode such as TE_1 , TE_2 , TE_3 like this.

Now, let us see how field varies in TE modes; so let us taken example of TE_2 mode.

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This is the electric field variation in TE 2 mode. This is the magnetic field variation in TE 2 mode. So, in the transverse electric mode there is only one component of electric field which is in y direction and there are two components of magnetic fields which are in X direction and Z direction H x and H z.

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And all of these 3 components are constant along y direction and they vary along X and Z only. So, we will see the variation of field in XZ plane only. So, let us first see the variation of electric field. So, this is X direction, this is Z direction and y direction is normal to this plane out word, and the dots in this circle represents the positive direction cross represent the negative Y direction and size of the circle represent the amplitude of the field.

So, E y is varying sinusoidally along X direction. So, if we put x is equal to 0 then it will be sin 0. So, we will get.

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0 field at x is equal to 0.

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If we put x is equal to a then we will get $\sin 2\pi$ which is also 0, so this will be 0 and if we put x is equal to a by 2 then $\sin \pi$ this will also be 0. So, electric field is 0 at x is equal to 0 x is equal to a by 2 and x is equal to a . And the maxima will be at x is equal to a by 4 and at x is equal to $3a$ by 4. So, there is a two half sinusoidal variation of electric field in x direction.

Now, let us see variation of electric field in Z direction. It will be 0 at Z is equal to 0 Z equal to λ by 2 Z is equal to λ and Z is equal to multiple of λ by 2, whereas it will be maximum at Z is equal to λ by 4 Z is equal to 3λ by 4 and Z is equal to odd multiple of λ by 4.

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So, this is how electric field varies in XZ plane for TE₂ mode. Now, let us see the variation of magnetic field in TE₂ mode. So, there are two magnetic fields and the resultant magnetic field will be vector sum of H_x and H_z . H_x is varying sinusoidally along x direction and H_z is varying co sinusoidally along X direction. It means these two components are in phase quadrature along X direction or we can say when H_x is maximum then H_z will be minimum or 0 and when H_x is 0 then H_z will be maximum which we can see from this figure also. At x is equal to 0 H_z is maximum and there is no field along x direction, so H_x is 0. At this point also x is equal to a by 2 H_z is maximum and H_x is 0 at this point also H_x is 0 and H_z is maximum.

So, H_z is maximum at x equal to 0 x equal to a by 2 and x equal to a , whereas it is minimum at x is equal to a by 4 and x is equal to $3a$ by 4. So, there is a two half sinusoidal variation of the fields along x direction.

Now, let us see the variation of these fields in a Z direction. These two components H_x and H_z are a step ahead by a length of λ by 4 because of this factor $j\beta$ there will be 90 degree phase delay. So, these two components are 90 degree out of phase. So, if H_x is maximum then H_z will be 0 along Z direction and if H_x is 0, then H_z will be maximum along Z direction.

So, if we will see in Z direction then H_z is maximum at Z is equal to 0 Z is equal to λ by 2 Z is equal to λ Z is equal to multiple of λ by 2, whereas H_x will

be 0 at these points Z is equal to 0, Z is equal to $\lambda/2$, Z is equal to λ like this.

Similarly, H_z will be 0 at Z is equal to $\lambda/4$, $3\lambda/4$ and odd multiple of $\lambda/4$ and at these points H_x will be a maxima. So, H_x is met maxima at Z is equal to $\lambda/4$, Z is equal to $3\lambda/4$ and odd multiple of $\lambda/4$. So, this is how electric and magnetic fields vary in TE₂ mode in parallel plane wave guide. So, in X direction or in the direction in which the wave is confined the variation of the fields is two half sinusoidal. Or in general if we have TEM mode or TMM mode then the variation of the fields along the direction in which the confinement has been done will be m half sinusoidal. This is all about the fields present in TM and TE modes.

In the next lecture we will discuss the cutoff frequencies for these modes, and after that we will start rectangular wave guides.

Thank you.