

Chemical Applications of Symmetry and Group Theory
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Lecture – 25

Hello, and welcome to the 25th class of this course. In the last class we learnt about the relation between Group Theory and Quantum Mechanics, and there we saw that the brief functions which are eigen functions of any given operator. For example, Hamiltonian operator which is an energy operator can act as the basis of for the irreducible representations for any particular molecule belonging into a particular point group.

Also we learnt about the direct products of the irreducible representations, which will be extremely useful in the coming weeks. And we also learnt how this direct product are used, they can be used. So, we took an example to show that direct product can be used to find out the energy elements, and also to find out the spectral intensity most precisely to probability of transition whether its allowed or disallowed and if it is allowed then, in which particular polarization it will be allowed.

Now one thing we did not talk about in the last class, that not only the wave function, which is an eigen function of an operator, can found the basis for irreducible representation, but also the linear combination of the wave function can also act as the basis for the same irreducible representation. Which is quite understood, because if I have an eigen functions, ψ_i , I for an operator, say H, the Hamiltonian then I can express ψ_i as a linear combination of various ϕ_i 's.

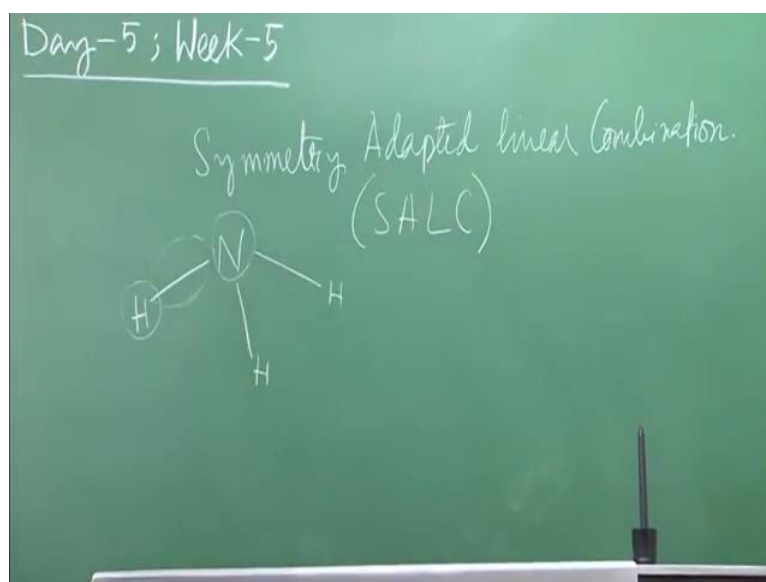
Therefore, its no wonder that proper linear combination, of the eigen functions can also act as a basis of the irreducible representation. And this property is of utmost importance when we try to find out which kind of say atomic orbitals will combine to give molecules orbitals, or which will contribute to the hybridized orbitals. Also to find out like which particular bond vectors will combine to keep an internal motion like vibration. So, in general the chemists they try to use the symmetry restrictions, and to facilitate themselves to understand chemical bonding and molecular dynamics. For example, as I said, constructing hybrid orbitals or molecular orbitals or finding proper orbitals sets under say ligand field or also analyzing vibrations of molecules.

In order to do that, that is developing such understanding based on symmetric restrictions one faces a common problem. The problem is that one needs to take one or more sets of orthonormal functions as required by a both group theory, and quanta mechanics. And here this orthonormal functions in case of this chemistry problems, this functions are taken to be the atomic orbitals. Or, in case of vibrational motions, are there other internal motions, they are internal coordinates of a molecules. In order to make proper combination, which will keep in the resultant function orthonormal, and that will act as a basis of irreducible representation one needs to be quite careful in his formulation of the problem.

And as we learnt in the last class this kind of the relation between the quanta mechanical observables, the wave functions rather not the observables and the symmetrical strength, this really is important for solving this problem.

So, in case when one tries to find out which atomic orbitals will form the hybrid molecular orbitals or in general molecular orbitals, one needs to take care of the symmetry constraints.

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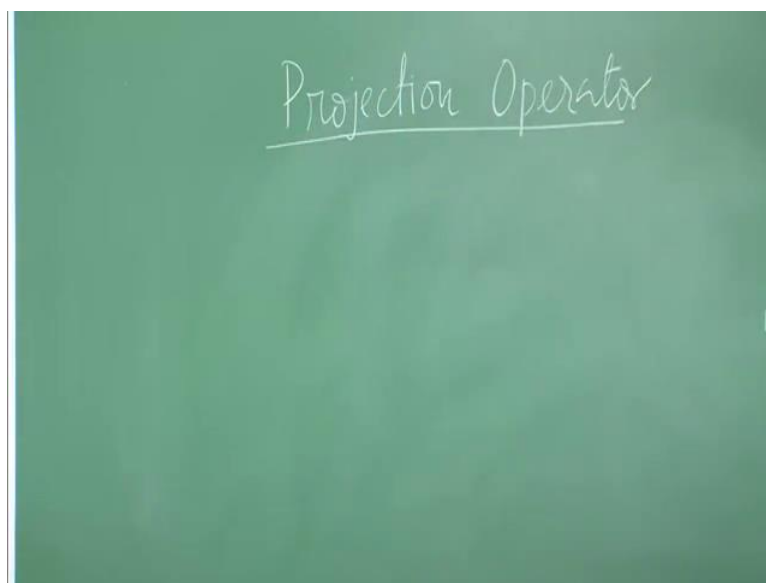


And then ultimately the linear combination that one gets, is known as Symmetry Adapted Linear Combination, or in short SALC. So, when SALC. So, how we can find out, that which orbitals will be able to combine in a linear fashion, and keeping the symmetric constraints to give a proper basis for a particular irreducible representation of particular

point group to which our constraints molecules belongs to. Now you see, why the symmetric constraints needed? Because for example, you take an any molecules; for example, ammonia. So, in the case of ammonia what we have is this bond that we write.

Now, bond means what? This is the you know combination of the atomic orbitals. So, hydrogen has one s orbitals, and nitrogen has 1s, 2s, 2 p orbitals. So, which orbitals will combine to form this successful bond, is not only governed by the energetics, but also governed by the symmetry. So, until unless the orbitals of this atoms, have the same symmetry they cannot combine. So, there is a symmetry restriction and therefore the you know the combination that we form to facilitate the bond formation, will be symmetry adapted linear combination, combination of atomic orbitals. So, how do you do that? So, there is a fundamental tool which is universally accepted and this particular tool is known as projection operator.

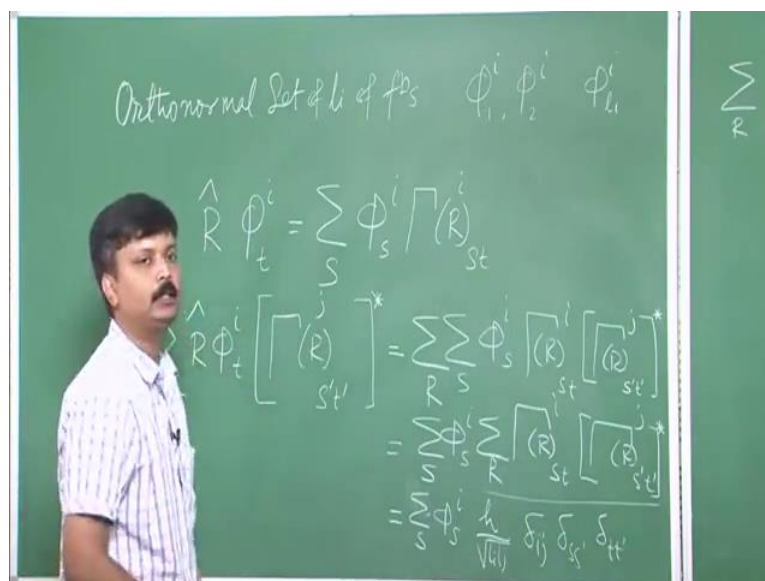
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So, now we will learnt about this particular projection operator, which will help us in finding on the SALCs, for any given molecule. In order to understand the position of at as functionally we better learn how to get this position operator first, and then will go for the actual application of this projection operator, by showing some illustrative examples. So, let us start with.

Let us assume that we have a orthonormal set of l_i number of functions, such as ϕ_1 ϕ_2 , up to ϕ_{l_i} .

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So, we have a l_i number of such functions, and we also assume that these functions form the basis of the i th irreducible representation of a point group, to which our concerned molecule belongs to. So, if this forms a basis for an i th irreducible representation, so we put a superscript as i , that will tell you that to which irreducible representation, these functions act as basis form. And the order of the group, that will be concerned, let us take as h , as usual. Now if I consider any operator, any symmetry operation as R , and I take any one of these functions, and operate any symmetry operation, on that function, then what should I get? So, if I take any arbitrary function ϕ_t , ϕ_t is 1 of these you know, l_i number of these, which formed of basis for i th irreducible representation.

So, that I can write as right? So, I can write this one, as I have taken this function to be the basis for i th irreducible representation, and this Γ_R having superscript i is a representation for this particular symmetry operation R , giving i th irreducible representation and st you can understand is a particular element. So, this Γ_R is a matrix representation. So, st gives me the particular element of that representation right? Now, what we have to do in order to, find out how to get the projection operator, what I will do I will multiply both the sides, with another representation for j th, i .

So, I will multiply both the sides, by a quantity which is Γ of the same operation R , and belongs to j th irreducible representation, and some element is s' , t' , and I will take a complex conjugate of that particular element, and after that we will sum it

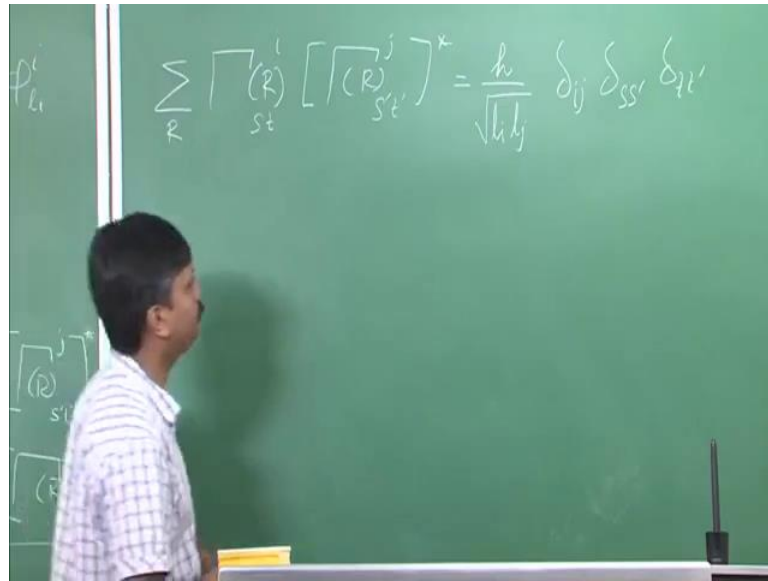
over all the symmetry operations r . So, we have to do the same thing on the right hand side as well. So, in that case I have summation over R , summation over s , $\phi_i s$ gamma of R , $s t$ multiplied by gamma of R , for j th irreducible representation having elements s prime, t prime ok?

So, now one thing you notice here, this ϕ is a function. So, it is independent of the symmetry operation right? It, it does not clash with the symmetry operation. So, I can you know, very easily separate this ϕ part, and this representation part. So, what I can do, I can rewrite this one as, summation over s ϕs i , and then summation over R and this part. That means, gamma R for i th irreducible representation, and element $s t$ multiplied by gamma of R , for j th irreducible representation having element s prime, t prime, and take the complex conjugate of this one. So, this we can very easily do right.

Now, we have a situation, where this right hand side tells me that I can have such l_i number of such products, where you know, each one will be a function ϕs that belongs that forms basis of (Refer Time: 25:05) irreducible representation and that is multiplied by a co-efficient. So, this part is a co-efficient which is multiplied this ϕs , and I can get such l_i number of products fine. Now, if I considered, this particular co-efficient, what does this co efficient tell me? So, this co efficient themselves R the summation of the products of 2 irreducible representations, over all the symmetry operations, that is what it tells me.

Now, such sum of the products of irreducible representation over all symmetry operations are governed by the great orthogonality theory. So, what does the great orthogonality theory tells me about such an such a co-efficient.

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So, if I just take this one. So, it is sum over R, gamma of R, for ith irreducible representation, and have been elements s and t, multiplied by gamma of R, for jth irreducible representation, having elements s prime, t prime, is equal to, if you remember the great orthogonality theorem, that we discussed earlier. So, we have considered the order of the group to be h. So, here order of the group, and then I have 2 different irreducible representations, ith and jth.

So their dimensions let us assume l I and l j, therefore, this will have l I and l j, right? and then I will have series of delta functions. So, here what we have? We have 2 irreducible representations get (Refer Time: 17:31) theorem will you know, guide you, how you can write the delta functions here. So, for 2 different irreducible representations, I will have the delta I j, and then, I have elements as s t for the ith, and s prime t prime for the jth irreducible representation therefore, I will have 2 more delta functions, one will be delta s s prime, and other 1 we will delta t, t prime alright? So, now what does that tell me? It tells me that I can replace this one over here.

So, this part will remain as it is, and I can write here, as h sorry, summation over s, phi s, for ith irreducible representation, and h by route over, l I, l j and then 3 delta functions, right? And then whatever I have on the other side that remains. So, here this right hand side immediately tells me that phi s, this side actually will give non 0 value, only when this is equals to s prime right? I am summing it over s. So, if I operate this operator on

this particular function ϕ_t , from i th representation, then I will have this right hand side having 1 0 value only when s is equals to s prime.

So, I will get $\phi_{s'}$ here, with 2 more delta function, δ_{ij} , $\delta_{t't}$. So, which means that not only that it has to be a particular function $\phi_{s'}$. In other to have non 0 value for this part, but also the 2 ire illusive representation that we are considering here must be identical meaning I equal to j and other element also should be the same that is t equals to t prime. Therefore, if consider all this what do I have is following.

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$$\sum_R [\Gamma_{s't'}^j]^* \hat{R} \phi_t^i = \frac{h}{h_j} \phi_{s'}^i \delta_{ij} \delta_{t't}$$

$$\frac{h_j}{h} \sum_R [\Gamma_{s't'}^j]^* \hat{R} \phi_t^i = \phi_{s'}^i \delta_{ij} \delta_{t't}$$

Projection operator $\hat{P}_{s't'}^j \phi_t^i = \phi_{s'}^i \delta_{ij} \delta_{t't}$

So, what we have here, the left side, sum over R , I can just slightly rearrange, because this 2 are not interfering each other. So, I will write this term first. So, I have γ of R , for j th ire illusive representation, having elements s prime t prime, multiplied by the operation R and optimum ϕ_t , forming basis for I , i th represent ire illusive representation. So, this is equals from here whatever I said if I can write that h by l_j , because when l_i , I equals to j , then only this right side I can write with survive. So, therefore, for I equals to j , I have l_i equals to l_j . So, l_j square and take a route, I get l_j . And I have ϕ , now s prime because for s equals to s prime, this will survive and for at this is representation. And if I still keep the deltas, took it the generality still, I will have this part.

Now, by rearranging, I can have $\int \phi_j^* \phi_i$ by $\int \delta_{ij}$, whatever we had, ϕ_j^* of s prime, $\int \delta_{ij}$, $\delta_{t t'}$, all right? So, this, this part of the left hand side, is known as the projection operator, and this is abbreviated as P of s prime, t prime, for j th irreducible representation. So, this is equal to the projection operators, and then you have this right? So, this is my projection operator. So, now, you can see here, that this particular operator when it is operated on any arbitrary function ϕ_t , it projects out another function $\phi_{s'}$.

So, this projection will take place, only when, this you know function ϕ_t , contains or it, it itself is $\phi_{s'}$. So, then it will give you know, that you know, survive, because I have this delta function $\delta_{t t'}$. Moreover it has to you know, ϕ_t , you know belong to the same particular irreducible representation. So, then I can have this $\phi_{s'}$, being projected out from an arbitrary function, or the combination of the function, which is given by ϕ_t .

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Orthogonal Set of li of fDs $\phi_1^i, \phi_2^i, \phi_l^i$

$$\hat{P}_{st}^j \phi_{t'}^j = \phi_{s'}^j$$

$$\hat{P}_{t't}^j \phi_t^i = \phi_{t'}^j \delta_{ij} \delta_{t't}$$

\hat{P}_{21}^j

So, ultimately by taking care of this delta functions, what I can right is this, making. So, I will rub this part. So, therefore, what I can write is, P of acting on general function, any arbitrary function t , which will give me ultimately $\phi_{t'}$, if I take care of this delta $t t'$.

So, essentially it means that this ϕ_t must contain some component of $\phi_{t'}$, in that case, I will have when $\phi_{t'}$, and for any arbitrary any irreducible

representation j , I can write this general form of this. Now this is the most general form, but I can have the special case. So, the special case is as follows. So, if I take this; so this is a particular special case, that we can consider. So, what does this special case tell me? So, it tells me that this P_j is the projection operator, which will project out ϕ_j , out of an arbitrary function which is ϕ that we started with right? So, by using such projection operators, based on the diagonal matrixes.

So, that is element rather we need generate from some arbitrary function, which is ϕ that we discussed, as you know a set of functions, which will form the basis of the j th irreducible representation. So, this is what we have as projection operator. Now this projection operator will be used in specific cases when we consider in the following class, taking a particular example of an molecule there we will consider the atomic orbitals there were functions and we will operate this projection operators here. So, here like, we say that this is particular irreducible representation j .

Now, what we can do we can you know take the real irreducible representation. So, you have to look at the character table find out the irreducible representation that you want to work with. So, for example, some irreducible representation b to u . So, can have that projection operator written as $P_{b \rightarrow u}$. So, in the following class we also see that there is, there are 2 type of projection operator one is like complete projection operator, another is incomplete projection operator.

So, most of our purpose we are pretty good with the incomplete projection operator itself which deals with not the matrix elements as such, but only with the characters. So, we can use the characters from character table belong to particular irreducible representation, to form the projection operator. And, thereby we can work on the particular orbitals of any you know atom within a molecule to project out which are the orbitals that can combine a linearly to form the you know say orbit orbitals.

In the following class we will learn more about this and try to form differently SALC's is from for different molecules; till then have a nice week.

Thank you very much.